We now turn to a particularly interesting type of systems of differential equations, known formally as a *gradient system*.

Definition. A *gradient system* on R^n is a system of differential equations of the form

$$
X' = -\text{grad}V(X) \tag{1}
$$

for $X \in \mathbb{R}^n$, where $V : \mathbb{R}^n \to \mathbb{R}$ is a smooth function with

$$
gradV = \left(\frac{\partial V}{\partial x_1}, \dots, \frac{\partial V}{\partial x_n}\right)
$$
\n(2)

The minus sign in [\(1\)](#page-0-0) is is merely a convention, and as such we need not worry of it's significance. As displayed in [\(2\)](#page-0-1), the gradient of a functions is simply an ordered list of it's partial derivatives. However, it's physical significance severely outweighs what it portrays to be be. The gradient of a point in a vector field is a measure of the *direction and rate of fastest increase*.

The following example is question 12, chapter 9 from Hirsch, Smale and Devaney's *Differential Equations, Dynamical Systems, and an Introduction to Chaos*.

Example 1.0. Let *T* be the torus defined as the square $0 \leq \theta_1, \theta_2 \leq 2\pi$ with opposite sides identified. Let $F(\theta_1, \theta_2) = \cos \theta_1 + \cos \theta_2$. Sketch the phase portrait for the system −grad*F* in *T*. Sketch a three-dimensional representation of this phase portrait with *T* represented as the surface of the doughnut.

Let $X =$ θ_1 θ_2 \setminus . Then we have the following system of equations;

$$
X' = \begin{pmatrix} \theta_1' \\ \theta_2' \end{pmatrix} = \begin{pmatrix} -\partial F/\partial \theta_1 \\ -\partial F/\partial \theta_2 \end{pmatrix} = \begin{pmatrix} \sin \theta_1 \\ \sin \theta_2 \end{pmatrix}
$$
 (3)

A first order, nonlinear uncoupled system. In order to understand the way this system behaves, we must solve for the equilibrium points of the system (namely, the pairs (θ_1, θ_2) such that $X' = 0$), and then linearize the system to understand the local behaviour near these points.

The reader may notice, however, that there are in fact an infinite number of pairs (θ_1, θ_2) which solve $X' = 0$: Any pair of the form $(n\pi, m\pi)$ for $n, m \in \mathbb{Q}$ will do. To make the task easier for us, we restrict our view to the square $[0, 2\pi]_{\theta_1} \times [0, 2\pi]_{\theta_2}$. In which case, we are left with 9 distinct equilibrium which can be summarized as $\{(n\pi, m\pi): n, m \in \{0, 1, 2\}\}.$ Now that we have our equilibrium points, we want to understand the local behaviour of solutions to [\(3\)](#page-0-2) in a neighbourhood of the equilibrium points. To do so, we take the Jacobian of *X*′ to get a linear planar system of the form $Y' = DF_XY$;

$$
DF_X = \begin{pmatrix} \cos \theta_1 & 0\\ 0 & \cos \theta_2 \end{pmatrix} \tag{4}
$$

If we consider the equilibrium point $(0,0)$, we find a linear system given by

$$
Y' = DF_{(0,0)}Y = \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} Y
$$
 (5)

A planar system with repeated, positive eigenvalues. This is indicates that (0*,* 0) is a source, and locally, solutions project directly outwards from this point. The identification of the other equilibrium are handled similarly, and are left to the reader. Figure [1](#page-1-0) shows the phase portrait of the system.

Figure 1: Phase portrait of [\(3\)](#page-0-2) on $[0, 2\pi] \times [0, 2\pi]$

For a more interactive view of this phase portrait, visit [Field Play](https://anvaka.github.io/fieldplay/?dt=0.01&fo=0.9999994&dp=0.01&cm=3&cx=3.12245&cy=2.9934499999999997&w=14.4621&h=14.4621&pc=30000&vf=vec2%20get_velocity%28vec2%20p%29%20%7B%0A%20%20vec2%20v%20%3D%20vec2%280.%2C%200.%29%3B%0A%0Av.x%20%3D%20sin%28p.x%29%3B%0Av.y%20%3D%20sin%28p.y%29%3B%0A%0A%20%20return%20v%3B%0A%7D&code=vec2%20get_velocity%28vec2%20p%29%20%7B%0A%20%20vec2%20v%20%3D%20vec2%280.%2C%200.%29%3B%0A%0Av.x%20%3D%20sin%28p.x%29%3B%0Av.y%20%3D%20sin%28p.y%29%3B%0A%0A%20%20return%20v%3B%0A%7D) \mathbb{Z} . As is shown, at each of the four corners, solutions flow out from a source, each of which falls into the sink located at (π, π) . Moreover, the midpoints of the perimeter each take the form of a saddle, with stable solutions flowing from the four corner, and the unstable lines again flowing directly towards the sink.

Moreover, we can *map* this phase portrait onto the torus. While we can certainly parameterize the surface of the torus using the map $T : [0, 2\pi]_u \times [0, 2\pi]_v \to \mathbb{R}^3$ defined by

$$
T(u, v) = [(2 + \cos u)\sin u, (2 + \cos uv)\cos u, \sin u]
$$
 (6)

It is, in my opinion, far more motivating (and interesting) to visualize the topological construction of the torus. More explicitly, we will construct the torus from a rectangular subset of \mathbb{R}^2 . See Appendix A??.

The phase portrait mapped onto the torus can be seen in Figure [2,](#page-2-0) with colour coordinated solutions which align with Figure [1.](#page-1-0) Only a sample of the lines have been plotted to allow for an easier viewing.

For a more visual view of this torus, visit [Math3D](https://www.math3d.org/VRcXdRUMt) \mathbb{Z} .